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The Crystal Structure of $\text{Hg}_3\text{O}_2\text{Cl}_2$, a Compound with Three-Coordinated Mercury(II)

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The structure of $\text{Hg}_3\text{O}_2\text{Cl}_2$ has been determined by Patterson and Fourier methods and refined to $R=0.059$ on the basis of 669 counter intensities. The crystals are monoclinic, space group $P2_1/c$, with $a=6.3100$ (7), $b=6.8657$ (5), $c=6.8579$ (5) Å, $\beta=114.366$ (6)° and $Z=2$. The mercury atoms are of two types, one [Hg(1)] linearly two-coordinated and the other [Hg(2)] three-coordinated with the mercury atom and its ligands in nearly the same plane. The mercury–oxygen distances are for Hg(1), 2.07 (2) Å ($2\times$) and for Hg(2), 2.17 (2), 2.32 (2) and 2.33 (2) Å. The shortest mercury–chlorine distance is 2.706 (5) Å. The mercury and oxygen atoms form a three-dimensional network with the chloride ions in its cavities. The oxygen and chlorine atoms are approximately hexagonally close-packed.

Introduction

The crystal chemistry of mercury oxide halides $n\text{HgO}\cdot\text{HgX}_2$ ($n=\frac{1}{2}, 2, 3, 4$) has been the subject of several investigations (Aurivillius, 1954, 1965). There seemed to be reasons for a re-investigation of the compound with $n=2$ (Aurivillius, 1954; Šćavničar, 1955), as the positions of the light atoms must be considered as uncertain and thus the mode of coordination of the mercury atoms not clearly revealed.

Work is also in progress on an additional compound with $n=2$ but with the formula $\text{Hg}_6\text{O}_4\text{Cl}_4$.

Experimental

Single crystals of $\text{Hg}_3\text{O}_2\text{Cl}_2$ were prepared by a method given by Arctowski (1895). An approximately 0.2 *M* aqueous solution of HgCl_2 was hydrolysed by the weak basic action of pieces of marble. After some weeks black crystals of irregular form had grown out from the surface of the marble. Larger quantities of the compound in microcrystalline form were synthesized by boiling a dilute aqueous solution of HgCl_2 with a deficit of mercury(II) oxide. After some minutes a black precipitate was obtained which gave Guinier–Hägg powder photographs identical with those obtained for the crystals synthesized according to Arctowski (1895). The sample thus prepared was analysed for mercury, chlorine and water, giving a result in good agreement with the formula $\text{Hg}_3\text{O}_2\text{Cl}_2$. Weissenberg photographs showed systematic absences $0k0$ for $k=2n+1$ and $h0l$ for $l=2n+1$, except for two weak $h0l$ reflexions. These two reflexions probably result from a small satellite on the crystal or a crystal defect. In spite of several attempts no other suitable single crystal could be picked out. Therefore the same crystal was used in the diffractometer work. The space group was assumed to be $P2_1/c$.

The density was determined from the loss of weight in benzene. Accurate cell dimensions were determined

from a least-squares analysis of the θ values of 39 reflexions measured on a single-crystal diffractometer.

Crystal data

$\text{Hg}_3\text{O}_2\text{Cl}_2$, F.W. 704.68.

Monoclinic, $P2_1/c$, $a=6.3100$ (7), $b=6.8657$ (5), $c=6.8579$ (5) Å,
 $\beta=114.366$ (6)°,
 $V=270.64$ Å³, $Z=2$.
 $D_m=8.53$, $D_x=8.59$ g cm⁻³. $\mu(\text{Mo } K\alpha)=869$ cm⁻¹.

Intensities were collected for the single crystal of approximate dimensions 0.10 × 0.10 × 0.14 mm on a computer-controlled diffractometer (CAD-4) with Mo $K\alpha$ ($\lambda=0.71069$ Å) radiation and a graphite monochromator.

The intensities were recorded at a take-off angle of 5°. The ω - 2θ scan technique was used with a scan interval of $(0.9+0.5 \tan \theta)^\circ$. The background was measured for $\frac{1}{4}$ of the scan time at each end of the interval. A fast pre-scan was used to calculate the scan speed at which a minimum net count of 2000 was attained within a maximum measuring time of 4 min. To test the systematic absences all reflexions with $3^\circ < \theta < 30^\circ$ were measured in two octants of reciprocal space. Of the 920 reflexions collected 70 were of type $h0l$ with $l=2n+1$ or $0k0$ with $k=2n+1$. Only four of these gave net intensities $I > 3\sigma(I)$, where $\sigma(I)$ is based on counting statistics. Two of them, $\bar{1}03$ and $\bar{3}09$, gave $I > 5\sigma(I)$ and were also observed in the Weissenberg photographs. These reflexions were rejected in the structure determination, assuming space group $P2_1/c$ as correct. The intensities of two standard reflexions, 004 and $\bar{3}31$, were measured at regular intervals to check the stability of the crystal and the electronics. No systematic variation was detected. All reflexions with $I < 3\sigma(I)$ were rejected. The remaining 669 independent reflexions were corrected for Lorentz, polarization and absorption effects. Transmission factors, evaluated by the numerical method, varied from 0.007 to 0.039.

Structure determination and refinement

The structure was solved by the heavy-atom method. From a vector map the six mercury atoms were found to occupy one fourfold and one twofold point position. A least-squares refinement of these positions, followed by difference syntheses, revealed the positions of the chlorine and oxygen atoms. A full-matrix least-squares calculation was then performed refining the atomic coordinates, isotropic temperature factors, and a scale factor. The resulting R was 0.10, where $R = \sum[|F_o| - |F_c|] / \sum|F_o|$. The function $\sum w_i(|F_o| - |F_c|)^2$ was minimized with the weights w_i calculated from the expression $w_i^{-1} = \sigma^2(F_o^2) / 4F_o^2 + aF_o^2 + b$. When anisotropic temperature factors were used for all atoms R fell to 0.07. A final refinement, including an isotropic extinction coefficient, converged to $R = 0.059$, $R_w = 0.073$ and $g = 0.26(3) \times 10^4$ where $R_w = [\sum w_i(|F_o| - |F_c|)^2 / \sum w_i |F_o|^2]^{1/2}$ and g is the extinction parameter.

In the last cycle all parameter shifts were less than 0.05 of the estimated standard deviations. The parameters a and b were 0.002 and 2.0, respectively. $S = [\sum w_i(|F_o| - |F_c|)^2 / (m - n)]^{1/2}$, where m and n are the number of observations and parameters varied, was 1.43. A final difference map was featureless.

The scattering factors were those of Doyle & Turner (1968), for neutral mercury, chlorine and oxygen. An anomalous dispersion correction for mercury (Cromer & Liberman, 1970) was also included.

The final positional parameters are given in Table 1 and the thermal parameters in Table 2. Observed and calculated structure amplitudes are presented in Table 3. In Table 4 selected interatomic distances and angles are given.

Table 1. Fractional coordinates obtained in the final least-squares refinement

Estimated standard deviations are given in parentheses.

$P2_1/c$; 2 Hg(1) in 2(a); 4 Hg(2), 4 Cl, 4 O in 4(e).

	x	y	z
Hg(1)	0	0	0
Hg(2)	0.4116 (1)	0.3635 (1)	0.2633 (1)
Cl	0.7891 (9)	0.1246 (8)	0.5546 (8)
O	0.3261 (25)	0.1084 (20)	0.0600 (22)

All computations were made on the UNIVAC 1108 computer in Lund. A short account of the program system used is given by Stålhandske (1974).

Description and discussion of the structure

Šćavničar (1955) has described the structure of this compound as being built up of Hg^{2+} ions and polymeric ions $(\text{OHgCl})_n^-$, the study being based on rather few intensities. The present investigation gives a result that differs from the earlier one.

The two unique mercury atoms Hg(1) and Hg(2) of the unit cell have different point symmetries, *viz.* Hg(1) $\bar{1}$ and Hg(2) 1. The atom Hg(1) is linearly coordinated to two oxygen atoms at 2.07(2) Å, whereas Hg(2) has a (1+2) coordination; to an oxygen atom at 2.17(2) Å and to two oxygen atoms at 2.32(2) and 2.33(2) Å. The linear coordination of Hg(1) is normal for a covalent *sp* bond of mercury, *e.g.* orthorhombic mercury(II) oxide (Aurivillius, 1956). The Hg(2)-O

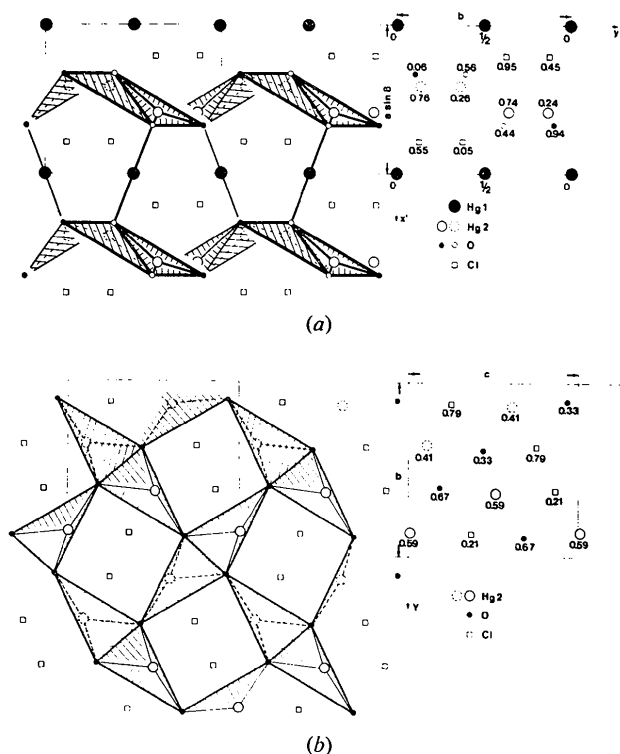


Fig. 1. (a) Projection of the structure along [001]. The oxygen atoms form the basal corners of the $\text{Hg}(2)\text{O}_3$ polyhedra, linked by Hg(1) atoms to form a three-dimensional network. The chloride ions are situated in cavities. (b) Projection of the structure on the yz plane. Hg(1) atoms are omitted.

Table 2. Vibrational parameters for the atoms

The anisotropic thermal parameters are based on the expression $\exp[-(h^2\beta_{11} + k^2\beta_{22} + l^2\beta_{33} + 2hk\beta_{12} + 2hl\beta_{13} + 2kl\beta_{23})]$. The β_{ij} values are multiplied by 10^5 and the r.m.s. components (Å) by 10^3 .

	β_{11}	β_{22}	β_{33}	β_{12}	β_{13}	β_{23}	R_1	R_2	R_3
Hg(1)	76 (3)	65 (2)	85 (3)	-12 (2)	31 (2)	0 (1)	106 (3)	128 (2)	134 (2)
Hg(2)	162 (3)	55 (2)	71 (2)	-3 (1)	29 (2)	-19 (1)	97 (2)	133 (2)	175 (2)
Cl	158 (14)	83 (10)	97 (11)	-5 (9)	50 (10)	0 (8)	138 (8)	141 (9)	165 (7)
O	137 (37)	36 (23)	112 (35)	0 (24)	75 (30)	13 (21)	87 (31)	126 (23)	162 (21)

distances are in the same range as in Hg(OH)₄Br₂, 2·25(5) Å (Aurivillius, 1968) for a tetrahedral mercury to oxygen connection. The next nearest neighbour

Table 4. Selected interatomic distances (Å) and angles (°) with standard deviations in parentheses

Table 3. Observed and calculated structure factor amplitudes

The columns are I, |F_o| and |F_c|.

Hg=Hg U	1 56 57	Hg=Hg 2	4 78 87	9 76 83	Hg=Hg 5	5 51 59
0 27 29	2 28 31	-4 55 59	6 184 203	6 58 62	-7 51 59	5 51 59
2 52 59	3 34 35	-3 60 65	6 61 66	8 49 50	-4 52 53	5 51 59
4 70 64	4 44 44	-7 67 69	6 38 38		-4 52 53	5 51 59
6 52 66	5 54 59	-1 62 63	9 92 92		-3 130 122	5 51 59
	6 35 37	0 121 106			-2 79 83	5 51 59
	6 26 26	1 121 119			-1 75 68	5 51 59
Hg=Hg 1					-1 75 68	5 51 59
0 41 44					-0 70 72	5 51 59
2 43 46					-3 11 16	5 51 59
3 93 87					-2 120 119	5 51 59
4 57 52					-3 26 29	5 51 59
6 44 41					0 167 164	5 51 59
7 78 73					2 145 150	5 51 59
					3 143 134	6 86 93
					4 109 127	
					6 112 121	
					7 87 96	
					8 64 69	
					9 81 88	
					10 113 108	
					11 151 151	
					12 144 145	
					13 23 22	
					14 83 83	
					15 117 115	
					16 129 112	
					17 30 32	
					18 59 62	
					19 54 57	
					20 63 64	
					21 40 36	
					22 31 35	
					23 40 41	
					24 27 29	
					25 37 35	
					26 68 67	
					27 45 47	
					28 44 41	
					29 33 32	
					30 173 174	
					31 183 186	
					32 25 23	
					33 46 45	
					34 23 25	
					35 51 51	
					36 48 45	
					37 75 74	
					38 40 38	
					39 57 55	
					40 21 25	
					41 35 30	
					42 118 111	
					43 25 23	
					44 173 174	
					45 28 27	
					46 47 44	
					47 19 13	
					48 28 23	
					49 118 111	
					50 25 23	
					51 173 174	
					52 28 27	
					53 46 45	
					54 23 25	
					55 51 51	
					56 48 45	
					57 75 74	
					58 40 38	
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					82 118 111	
					83 25 23	
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					230 25 23	
					231 173 174	
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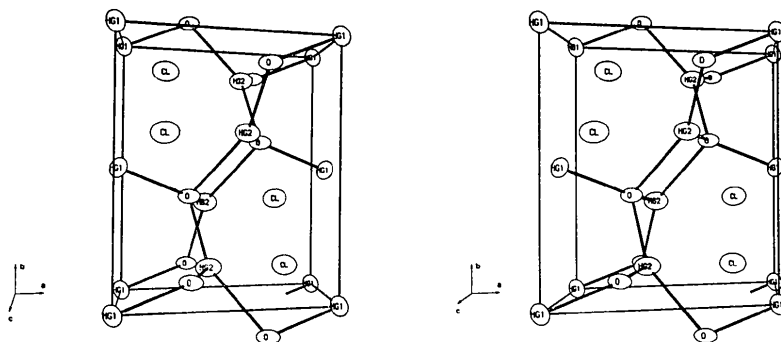


Fig. 2. A stereo view of the structure.

to three mercury atoms Hg(2) at 2.17, 2.32 and 2.33 Å. Each polyhedron thus shares its basal corners with three polyhedra and its basal edges with two polyhedra, building up infinite layers parallel to (100). Every oxygen atom is also bonded to an atom Hg(1) at 2.07 Å, completing the deformed tetrahedral coordination of oxygen. Thus every basal corner of the polyhedra in parallel layers is linked by the atoms Hg(1) to form a three-dimensional network of formula $(\text{Hg}_3\text{O}_2)_n^{2n+}$.

The mercury–chlorine distances are $\geq 2.706(5)$ Å (Table 4), indicating mainly ionic bonding. The chloride ions are situated in cavities in the three-dimensional network. A stereo view of the structure is given in Fig. 2.

From an *entirely geometrical* point of view the polyhedra around Hg(1) and Hg(2) can be described as follows:

Each Hg(1) atom is surrounded by two oxygen atoms at short distances (2.07 Å) and four chlorine atoms more distant (2.92–3.00 Å), the atoms forming a distorted octahedron with the chlorine atoms defining an equatorial plane. Hg(2) atoms are each surrounded by three oxygen atoms (2.17–2.33 Å) and a chlorine atom (2.71 Å), the atoms forming a deformed tetrahedron with the mercury atom drawn very near to one face. If the next nearest chlorine atom at 2.91 Å is also considered, the polyhedron is a deformed trigonal bipyramid.

The arrangement of the oxygen and chlorine atoms approximates hexagonal close-packing.

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